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# 4f hybridization and band dispersion in gadolinium thin films and compounds

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There is interplay between intra-atomic orbital hybridization and extra-atomic hybridization in various gadolinium systems, which affects magnetic coupling and electron itinerancy (localization). The results do not always follow expectation. The experimental band structure of thin Gd(0001) films, grown on the Mo(112) surface, along the  $\bar{\Gamma}-\bar{M}$  does not agree with ex-

pectations even qualitatively. In particular, the dispersion of the gadolinium band, with strong 5d weight near 2 eV binding energy provides considerable evidence to support the case for 4f-5d hybridization, with increasing 5d localization. On the other hand, there is also evidence of extra-atomic Gd 4f hybridization leading band dispersion in the occupied 4f levels in Gd<sub>2</sub>O<sub>3</sub>.

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In local moment 4f systems, two issues that have a strong influence on our understanding of exchange coupling are the extent of 4f localization or itinerancy AND the extent of 4f–5d hybridization. While a number of experimental studies have been undertaken to investigate the spin-polarized band structure of strained Gd(0001) [1–4], interpretation of experimental data is complicated by the possibility of 4f–5d hybridization. Singh [5] has suggested that there is hybridization with the upper part of the gadolinium 4f manifold (spin minority 4f) with the unoccupied 5d/6s bands, and this is very consistent with the experimental band structure [4]. Here we address the issue of 4f–5d hybridization with respect to the spin majority 4f<sup>7</sup> gadolinium (lower part) manifold.

Without a treatment of the 4f electrons that includes hybridization, correlation energies, and band structure, the calculations (despite the diversity of approaches) have exhibited only limited agreement with photoemission experiments [4]. In the case of strained Gd(0001), comparisons of the occupied band structure with theory are dramatically worse than for strain relieved Gd(0001) grown on W(110) [4, 6, 7].

It has been established that the local spin density approximation (LSDA) fails to describe correctly electronic structure of strongly correlated systems such as the 4f metals [5, 8, 9]. The inability of LSDA to correctly describe 4f metals appears to occur mainly because of the incorrect position of spin up  $f\uparrow$  and spin down  $f\downarrow$  bands if the 4f elec-

trons are treated as bands and not as core states [10–12]. Harmon and Freeman [13] calculated the band structure for Gd by APW method using the exchange-correlation  $X_\alpha$  potential with two values i.e.  $\alpha = 1$  (“total Slater’s exchange”) and  $\alpha = 2/3$  (“Kohn-Sham potential”). The occupied  $4f^7$  states were treated as core states and given a rigid value of magnetic moment  $\mu = 7\mu_B$  associated with spin-up  $4f^7$  electrons. In order to get the experimental value of  $7.55 \mu_B$  for Gd, the rest magnetic moment, some  $0.55 \mu_B$ , was provided by properly choosing the configuration of the 5d6s valence electrons. The results for  $\alpha = 1$  were in reasonable agreement with experiment, though the obtained width of the valence band (from the bottom at  $\Gamma$  to  $E_F$ ) of 3.1 eV seems too small compared to the values derived from photoemission of about 4 eV.

Harmon et al. [9] and Singh [5] suggested that correct total energies for Gd could not be obtained if 4f electrons were treated as core states. The increasingly compelling evidence that gadolinium was a local moment system, with correlated electrons [7], provoked the development of an improved version of LSDA by including the Hubbard U, taken from the mean-field approach [8]. The LSDA + U method has proved fruitful in recent calculations of the bulk and surface electronic structure of relatively unstrained Gd [14–16]. Due to a better accounting of the intra-atomic correlations, the 4f minority band rises away from Fermi level while the 4f majority band increases in binding energy to the correct position [14] (approximately  $-8.5$  eV with respect to  $E_F$  [7]). Wu et al. [17], using von Barth and Hedin’s exchange-correlation potential, calculated the Gd(0001) band structure by the FLAPW method. The majority  $4f^7$  electrons were treated as core states. Despite using an occupied s-d band width (almost 6 eV) that is too large, the results are in qualitative agreement with angle-resolved photoemission [18–20] and inverse photoemission [21–23] data, again of the relatively unstrained Gd(0001). In particular, the calculations reveal the existence of the spin majority (in vicinity of  $\bar{\Gamma}$ ) surface band of  $d_{3z^2-r^2}$  symmetry 0.3 eV below  $E_F$  [4,17,19] as well as spin minority surface resonance bands above  $E_F$  which are observed in the spin-polarized photoemission [2,24] and spin-polarized inverse photoemission [22,23] spectra respectively. The issue of correct description of 4f states was also addressed by Eriksson et al. [12] where it was suggested that the “first-principle” (pure) LSDA approach could be resuscitated, provided that the 4f states are treated as core states.

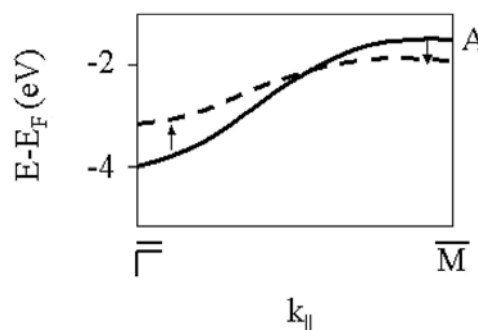
Following Bylander and Kleinman [10], Yakovkin et al. [4] adopted the parameter description for correlation energy, while the core electrons are treated within the same approximation for exchange-correlation potential.

With expansive strain, increased localization will increase the 5d band localization and might well result in greater *intra-atomic* overlap between 4f and 5d levels [6]. This, in-turn, should also affect the bottom of the 5d/6s bulk bands as indicated by angle resolved photoemission

band mapping, reported elsewhere [4,6]. Reasonable agreement between experiment and theory was realized excepting the highest binding energy 5d/6s Gd bulk band - the occupied bulk band most likely to provide indications of occupied 4f-5d hybridization. If this is indeed the case, then it is questionable whether the Gd 4f levels can be treated as pure core states. If the 4f levels cannot be treated as shallow core level states, then a proper treatment of the 4f levels needs to be revisited.

**The possible effect of strain on Gd 5d localization:** The expansive 4% in-plane strained Gd(0001) on Mo(112) results in a quite different electronic structure and altered magnetic properties compared to the strain relieved Gd(0001) grown on W(110) [25], as has been noted elsewhere [2,4,6,26]. Based on studies of rare earth alloys, Andrianov [27–29] suggested that the Fermi surface is sensitive to both composition and lattice deformation. Gd(0001) grown on Mo(112) exhibit expansive strain and an increased in-plane lattice constant, expanded by 4% for a film thickness of 30 to 150 Å [6]. Several interesting characteristics of such strained Gd films have already been noted: large magneto-elastic interactions [30,31], enhance Curie temperatures [1], a different electronic band structure compared to that of bulk “unstrained” Gd [2–6], and others as described below. Such experimental observations have also been noted for other rare earths as well [34–36].

In comparing the calculated electronic band structures of bulk “unstrained” Gd and the strained Gd film, two significant differences emerge in both the calculations of Yakovkin [4] and Freeman [37] as compared with experiment. As schematically shown in Fig. 1, the largest binding energy 5d/6s bulk band exhibits less dispersion with increasing expansive strain. The Gd slab calculations result in narrower occupied band (smaller dispersion) with increasing lattice constant, yet the direction of the dispersion remains the same. In addition, regardless of strain, the band is  $\Delta_1$  or  $\alpha_1$  symmetry ( $5d_{z^2-r^2}$ , 6s character), even with extensive expansive strain.



**Figure 1** This is a schematic band structure showing the dispersions of the highest binding energy 5d bulk bands. Solid line represents the unstrained Gd (0001) and dashed line is for the expansively strained Gd.

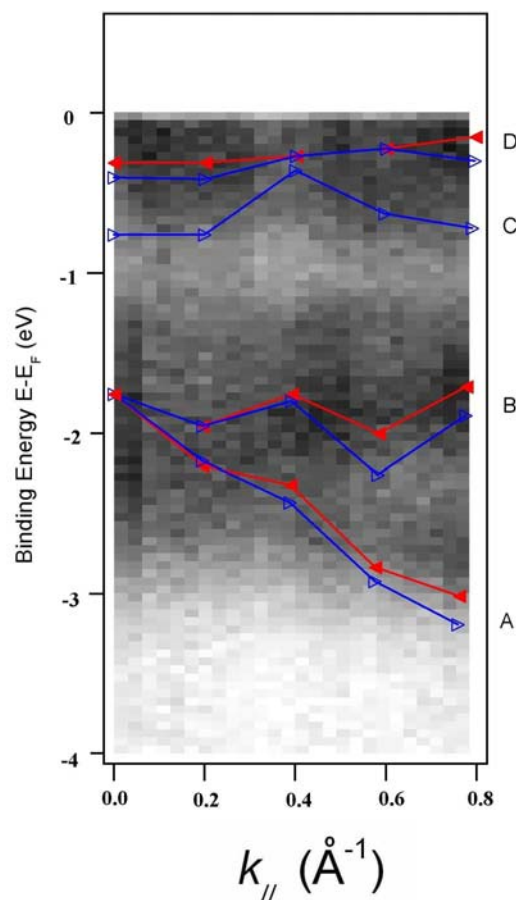
These expectations from theory are not reproduced in experiment [6]. These bulk bands are experimentally observed to change symmetry from  $\Delta_1, \Delta_2$  ( $5d_{z^2-r^2}$ , 6s) for unstrained Gd(0001) on W(110) to  $\Delta_5, \Delta_6$  at  $\bar{\Gamma}$  with the 4% expansive strain [6]. The second difference between experiment and theory is the band dispersion of the highest binding energy bulk bands, discussed herein.

We have mapped out band structure of thin Gd films grown on the Mo (112) surface along the  $\bar{\Gamma}-\bar{M}$  high-symmetry direction of the surface Brillouin zone. The experiments were carried out in an UHV chamber equipped with a Scienta 200 electron analyzer, at the Synchrotron Radiation Center in Stoughton, Wisconsin, with considerable wave vector and energy resolution in angle-resolved photoemission. The clean Mo (112) surface was obtained after the standard surface preparation treatment [6], which included heating in oxygen atmosphere with subsequent serial flashing to the temperatures above 2100 K. The Gd (0001) films were grown at room temperature at the base pressure of  $1 \times 10^{-10}$  Torr and subsequently annealed. The crystal quality of the Gd (0001) films was determined by LEED, which was also used to determine the extent of the expansive strain, along with the critical points in the experimental band mapping. Spectra from the photoemission are also used to confirm crystal quality since contaminated spectra are definitely different [6, 7]. The in-plane lattice constant is expanded by 4% compared to that of the bulk Gd ( $a = 3.6$  Å).

Angle-resolved spectra were taken at around 200 K and the spin-integrated wave vector spectral density is plotted in Fig. 2. The spin integrated data can be easily compared with the spin polarized band mapping [2–4]. The slight displacement in the wave vector dependent binding energies between spin-polarized data and the spin integrated data are due to changes in  $k_{\perp}$  (the spin integrated spectral density data were taken at 32 eV while the spin resolved band mapping was acquired at 35 eV) affecting principally the bulk bands. In spite of the difference in photon energy between the two sets of data, the spin majority (blue) and spin minority (red) bands agree relatively well with the wave vector dependent spectral density from present angle-resolved and spin-integrated spectra.

While the bands around the Fermi energy are described very well elsewhere [2–4], a repetition of this is beyond the scope of the present paper. The bulk 5d6s bands (labeled A and B in Fig. 2) do not show the expected positive slope in the dispersion compared to unstrained Gd due to the strain and the second bulk band B, nicely resolved away from  $\bar{\Gamma}$ , is not expected from theory at all. This band is a result of strain but the origin of this is more difficult to explain.

Theory, in which the 4f states are treated as core states or a perturbation to the valence band structure could be deficient. The bulk bands we label the 'A' and 'B' bands may arise from combination of 4f and 5d wave functions. We would expect more weight of 5d states near  $\bar{\Gamma}$  and more 4f states near  $\bar{M}$ . This profound disagreement between the-



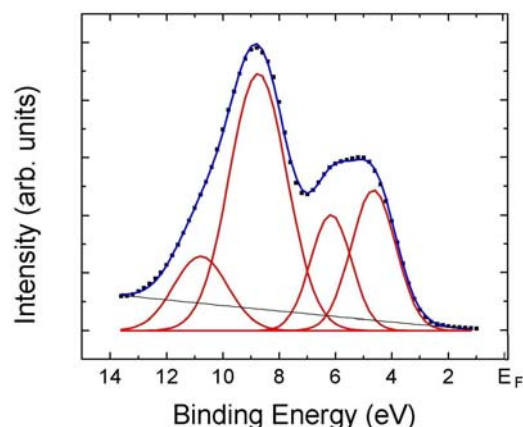
**Figure 2** Spin-integrated Scienta data at 32 eV, with an overlay of the band dispersion inferred from spin resolved photoemission taken at a photon energy of 35 eV, and select wave vectors. The latter data is adapted from reference [2,3]. In the spin resolved data overlay, red triangles are spin minority and blue is spin majority. Higher photoemission intensities are shown dark.

ory and experiments we take to be an indication of 4f–5d mixing.

While considerable agreement with experiment was obtained [4], there are two significant differences between theory and the experimental band mappings both in regard to the strained Gd films the 5d, 6s occupied bulk spin majority and minority subbands found at approximately 1.8 eV binding energy (at 35 eV photon energy) to 2.0 eV binding (at 32 eV photon energy) at  $\bar{\Gamma}$ . Light polarization photoemission experiments indicate that these bulk bands change symmetry from  $\Delta_1, \Delta_2$  ( $5d_{z^2-r^2}$ , 6s) for unstrained Gd(0001) on W(110) to  $\Delta_5, \Delta_6$  at  $\bar{\Gamma}$  with the 4% expansive strain for Gd(0001) grown on Mo(112)[6]. No matter how extensive the strain, expansive strain does not result in a change in symmetry of these bands from  $\Delta_1, \Delta_2$  ( $5d_{z^2-r^2}$ , 6s).

Since strained Gd film has a greater lattice constant, owing to the slight expansive strain (4%) compared to Gd





**Figure 3** The photoemission spectrum of the valence band for  $\text{Gd}_2\text{O}_3(-402)$ . The various components contributing to the valence band structure are indicated. The photon energy is 100 eV and the light incidence angle is  $45^\circ$ . All photoelectrons were collected along the surface normal at  $T = 240^\circ\text{C}$ . Binding energies are in  $E_F - E$ .

bulk, adjacent Gd atomic wave functions decrease in the extent of overlap. This, in turn, could increase the localization of the Gd 5d band and thus increase the 5d–4f wave function overlap and the potential for intra-atomic hybridization. 5d band overlap may be a key to 4f/5d hybridization. In the more localized rare earth compounds, Olson et al. [32] found evidence for 4f delocalization in  $\text{CeSb}_2$ , in the form of dispersing Bloch states. Of course f delocalization can also occur, but is more expected for the 5f levels. Joyce et al. [33] also found clear evidence of similar dispersion in the uranium 5f levels. In some sense, the result here are similar to the apparent effects of 4f–5d hybridization suggested by the result seen for  $\text{ErAs}(100)$ , where strong 4f intra-atomic hybridization was implicated, and a direct confirmation of an occupied 4f band structure inferred from the data but could not be established [38,39].

**Gd 4f band structure:** In the  $\text{Gd}^{3+}$  and mixed valence systems, Gd 4f hybridization with nearest neighbor atoms is expected [40,41] and observed. To actually demonstrate this, we have mapped out the bulk band structure of the Gd 4f weighted bands in a system where the Gd 5d/6s bands are heavily suppressed:  $\text{Gd}_2\text{O}_3$ . This experimental band mapping was done along the normal to  $(-402)$ , using angle-resolved photoemission. Here, experiments were performed using the 3m toroidal grating monochromator (3 m TGM) beam line [42] in a UHV chamber previously described [38,39,42]. The Fermi level ( $E_F$ ) was established from a gold film in electrical contact with the sample and measurements were carried out at ambient temperatures.

When  $\text{Gd}_2\text{O}_3$  is grown on Si(100) using a pure  $\text{Gd}_2\text{O}_3$  target by pulsed laser deposition, the substrate has a strong

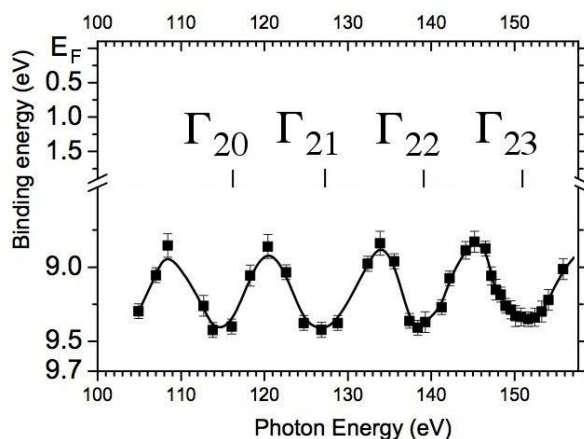
influence on the texture and crystal structure of the  $\text{Gd}_2\text{O}_3$  film [43]. Such films are highly textured monoclinic  $\text{Gd}_2\text{O}_3$ , not the more expected cubic structure [44–46]. The textured structure is such that the  $(-402)$  planes mostly lie along the surface of the film [43]. The lattice spacing along the  $\langle -402 \rangle$  direction is about 0.2965 nm so the repeat along the surface normal is eight times the  $(-402)$  layer spacing or about 2.372 nm. So while the  $\text{Gd}_2\text{O}_3$  unit cell is large with inequivalent Gd (three) and oxygen atoms (in a number of inequivalent sites), it is a system in which some 4f band dispersion might be observed, in spite of some 84 gadolinium bands per unit cell.

The valence band density of states, as determined by photoemission for  $\text{Gd}_2\text{O}_3$ , are very similar to that observed for 3% Gd-doped hafnium oxide films [41,42,47] and reported elsewhere for  $\text{Gd}_2\text{O}_3$  [48,49], as indicated in Fig. 3. The major difference is that for  $\text{Gd}_2\text{O}_3$ , we require 4 components at a minimum to fit the valence band spectra, as also indicated in Fig. 3, as opposed to three components needed for Gd doped  $\text{HfO}_2$ , in the monoclinic phase [41,42,47]. Confirmation that the Gd 4f states contribute to the photoemission components at about 8.7 to 9.5 eV and the ‘shoulder’ at 10 to 11 eV binding energies is found in resonant photoemission (i.e., constant initial state spectroscopy) measurements of Gd doped  $\text{HfO}_2$  [42,47]. The photoelectron intensities from Gd doped  $\text{HfO}_2$ , determined from the feature at about 9.5 eV binding energy (from the Fermi level) is strongly enhanced at about 152 eV, or at the  $4d^{10}4f^7 \rightarrow 4d^9 4f^8 \rightarrow 4d^{10} 4f^6 + e^-$  super Coster–Kronig transition [40,41,41,50,51], leading to a classic Fano resonance.

Since the two components, identified as heavily Gd 4f in weight, are separated by slightly more than 2 eV, we must conclude that the higher binding energy component is, in fact, a consequence of the surface to bulk shift, and is the surface component. The Gd occupied 4f surface-to-bulk shift is only about 0.4 eV for Gd metal, but is expected to be much larger with oxidation [49,52], with the surface component at the greater binding energy. It is the larger intensity component, at 8.7 to 9.5 eV binding energy, associated with the bulk Gd 4f weighted bands of  $\text{Gd}_2\text{O}_3$ , that are of interest in establishing the Gd 4f band structure.

Because of the high degree of order and the strong texture growth along  $\langle -402 \rangle$ , we have been able to plot out the bulk band structure along  $\langle -402 \rangle$ . Because the photoelectrons are collected along the surface normal in the sequence of photon energy dependent photoemission spectra, the binding energy shifts of many of the observed photoemission features is indicative of band dispersion along the electron wave vector normal to the surface,  $k_\perp$ . The value of  $k_\perp$  can be estimated from the photoelectron kinetic energy making some assumptions about the inner potential  $U_{\text{in}}$ :

$$k_\perp = \sqrt{\frac{2m}{\hbar^2} \{E_{\text{kin}}(\cos(\theta))^2 + U_{\text{in}}\}}. \quad (1)$$



**Figure 4** The dispersion of the Gd 4f component of the valence band of gadolinium oxide, with changing photon energy. The critical points are indicated, assuming no inner potential and the predicted lattice spacing of 23.7 Å. Binding energies are in  $E_F - E$ .

The dispersion of the component we attribute to the Gd 4f dependence level at 9 to 9.5 eV binding energies is summarized in Fig. 4. The dependence on photon energy shows band critical points repetition suggesting periodicity of about  $22 \pm 2$  Å perpendicular to the film or along the surface normal, close to the value of 23 Å as expected from the crystallography.

For a dielectric insulator we would expect that the value of the inner potential to be effectively negligible, as seems to be the case here. We cannot establish that the inner potential is zero, but from our measurements that establishes the critical points for the 19th through to the 23rd Brillouin zone, the inner potential is quite small and is no more than 1–2 eV, at most. Assuming a value of zero for the inner potential, and a lattice constant of 23.72 Å, the calculated critical points are in good agreement with the experimental band structure, as indicated in Fig. 4.

What seems clear is that not only do we have strong Gd 4f hybridization with the oxygen for Gd<sub>2</sub>O<sub>3</sub> and we have a small amount of band dispersion in the band strongly weighted with a Gd 4f contribution for Gd<sub>2</sub>O<sub>3</sub>. The band dispersion is small (about 500 meV), but evident nonetheless. The Gd 4f weighted band is placed well away from the Fermi level in the valence band of Gd<sub>2</sub>O<sub>3</sub> but still exhibits band structure.

**Summary:** We have mapped out band structure of thin Gd films grown on the Mo (112) surface along the  $\bar{\Gamma}-\bar{M}$  and the results are largely consistent with previous work [1–4]. Nonetheless, strong deviations from theory are reproduced, and this in turn implicates 4f–5d hybridization. This 4f–5d hybridization appear now not only to occur for the unoccupied band structure [4,5] but may exist for the

occupied Gd 4f band structure as well. The hybridization for the unoccupied band structure results from the energy position (overlap) of the d–f bands, which is not so very evident for the occupied bands. In spite of the high index face adopted by Gd<sub>2</sub>O<sub>3</sub> thin films on Si(100), and the large unit cell of Gd<sub>2</sub>O<sub>3</sub>, Gd 4f band structure is nonetheless and perhaps surprisingly apparent for Gd<sub>2</sub>O<sub>3</sub>.

This raises some interesting issues that have not yet been successfully addressed in theory. First, how can we account for the observed change in symmetry from  $\Delta_1, \Delta_2$  ( $5d_{z^2-r^2}, 6s$ ) for unstrained Gd(0001) on W(110) to  $\Delta_5, \Delta_6$  at  $\bar{\Gamma}$  [6] and the resulting influence of the slope of the band dispersion. Further, does this change in valence band symmetry indeed imply intra-atomic Gd 4f–5d hybridization? Second, are the Gd 4f levels really shallow core levels, or states with a weak band dispersion?

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